

## 3.9 Fluid dynamics

### Ideal fluids<sup>a</sup>

Continuity <sup>b</sup>	$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$	(3.285)	$\rho$ density
Kelvin circulation	$\Gamma = \oint \mathbf{v} \cdot d\mathbf{l} = \text{constant}$	(3.286)	$\mathbf{v}$ fluid velocity field
	$= \int_S \boldsymbol{\omega} \cdot d\mathbf{s}$	(3.287)	$t$ time
Euler's equation <sup>c</sup>	$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = -\frac{\nabla p}{\rho} + \mathbf{g}$	(3.288)	$\Gamma$ circulation
	or $\frac{\partial}{\partial t}(\nabla \times \mathbf{v}) = \nabla \times [\mathbf{v} \times (\nabla \times \mathbf{v})]$	(3.289)	$d\mathbf{l}$ loop element
Bernoulli's equation (incompressible flow)	$\frac{1}{2} \rho v^2 + p + \rho g z = \text{constant}$	(3.290)	$ds$ element of surface bounded by loop
Bernoulli's equation (compressible adiabatic flow) <sup>d</sup>	$\frac{1}{2} v^2 + \frac{\gamma}{\gamma-1} \frac{p}{\rho} + g z = \text{constant}$	(3.291)	$\boldsymbol{\omega}$ vorticity ( $= \nabla \times \mathbf{v}$ )
	$= \frac{1}{2} v^2 + c_p T + g z$	(3.292)	$p$ pressure
Hydrostatics	$\nabla p = \rho \mathbf{g}$	(3.293)	$\mathbf{g}$ gravitational field strength
Adiabatic lapse rate (ideal gas)	$\frac{dT}{dz} = -\frac{g}{c_p}$	(3.294)	$(\mathbf{v} \cdot \nabla)$ advective operator
<i>a</i> No thermal conductivity or viscosity.			
<i>b</i> True generally.			
<i>c</i> The second form of Euler's equation applies to incompressible flow only.			
<i>d</i> Equation (3.292) is true only for an ideal gas.			

### Potential flow<sup>a</sup>

Velocity potential	$\mathbf{v} = \nabla \phi$	(3.295)	$\mathbf{v}$ velocity
	$\nabla^2 \phi = 0$	(3.296)	$\phi$ velocity potential
Vorticity condition	$\boldsymbol{\omega} = \nabla \times \mathbf{v} = 0$	(3.297)	$\boldsymbol{\omega}$ vorticity
Drag force on a sphere <sup>b</sup>	$\mathbf{F} = -\frac{2}{3} \pi \rho a^3 \dot{\mathbf{u}} = -\frac{1}{2} M_d \ddot{\mathbf{u}}$	(3.298)	$F$ drag force on moving sphere
<i>a</i> For incompressible fluids.			
<i>b</i> The effect of this drag force is to give the sphere an additional effective mass equal to half the mass of fluid displaced.			
$a$ sphere radius			
$\dot{\mathbf{u}}$ sphere acceleration			
$\rho$ fluid density			
$M_d$ displaced fluid mass			

## Viscous flow (incompressible)<sup>a</sup>

Fluid stress	$\tau_{ij} = -p\delta_{ij} + \eta \left( \frac{\partial v_i}{\partial x_j} + \frac{\partial v_j}{\partial x_i} \right)$	(3.299)	$\tau_{ij}$ fluid stress tensor $p$ hydrostatic pressure $\eta$ shear viscosity $v_i$ velocity along $i$ axis $\delta_{ij}$ Kronecker delta
Navier–Stokes equation <sup>b</sup>	$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = -\frac{\nabla p}{\rho} - \frac{\eta}{\rho} \nabla \times \boldsymbol{\omega} + \mathbf{g}$	(3.300)	$\mathbf{v}$ fluid velocity field $\boldsymbol{\omega}$ vorticity $\mathbf{g}$ gravitational acceleration
Kinematic viscosity	$v = \eta / \rho$	(3.302)	$\rho$ density $v$ kinematic viscosity

<sup>a</sup>I.e.,  $\nabla \cdot \mathbf{v} = 0$ ,  $\eta \neq 0$ .

<sup>b</sup>Neglecting bulk (second) viscosity.

## Laminar viscous flow

Between parallel plates	$v_z(y) = \frac{1}{2\eta} y(h-y) \frac{\partial p}{\partial z}$	(3.303)	$v_z$ flow velocity $z$ direction of flow $y$ distance from plate $\eta$ shear viscosity $p$ pressure
Along a circular pipe <sup>a</sup>	$v_z(r) = \frac{1}{4\eta} (a^2 - r^2) \frac{\partial p}{\partial z}$	(3.304)	$r$ distance from pipe axis $a$ pipe radius
	$Q = \frac{dV}{dt} = \frac{\pi a^4}{8\eta} \frac{\partial p}{\partial z}$	(3.305)	$V$ volume
Circulating between concentric rotating cylinders <sup>b</sup>	$G_z = \frac{4\pi\eta a_1^2 a_2^2}{a_2^2 - a_1^2} (\omega_2 - \omega_1)$	(3.306)	$G_z$ axial couple between cylinders per unit length $\omega_i$ angular velocity of $i$ th cylinder
Along an annular pipe	$Q = \frac{\pi}{8\eta} \frac{\partial p}{\partial z} \left[ a_2^4 - a_1^4 - \frac{(a_2^2 - a_1^2)^2}{\ln(a_2/a_1)} \right]$	(3.307)	$a_1$ inner radius $a_2$ outer radius $Q$ volume discharge rate

<sup>a</sup>Poiseuille flow.

<sup>b</sup>Couette flow.

## Drag<sup>a</sup>

On a sphere (Stokes's law)	$F = 6\pi a \eta v$	(3.308)	$F$ drag force $a$ radius
On a disk, broadside to flow	$F = 16a\eta v$	(3.309)	$v$ velocity
On a disk, edge on to flow	$F = 32a\eta v/3$	(3.310)	$\eta$ shear viscosity

<sup>a</sup>For Reynolds numbers  $\ll 1$ .

## Characteristic numbers

Reynolds number	$\text{Re} = \frac{\rho UL}{\eta} = \frac{\text{inertial force}}{\text{viscous force}}$	(3.311)	$\text{Re}$ Reynolds number $\rho$ density $U$ characteristic velocity $L$ characteristic scale-length $\eta$ shear viscosity
Froude number <sup>a</sup>	$F = \frac{U^2}{Lg} = \frac{\text{inertial force}}{\text{gravitational force}}$	(3.312)	$F$ Froude number $g$ gravitational acceleration
Strouhal number <sup>b</sup>	$S = \frac{U\tau}{L} = \frac{\text{evolution scale}}{\text{physical scale}}$	(3.313)	$S$ Strouhal number $\tau$ characteristic timescale
Prandtl number	$P = \frac{\eta c_p}{\lambda} = \frac{\text{momentum transport}}{\text{heat transport}}$	(3.314)	$P$ Prandtl number $c_p$ Specific heat capacity at constant pressure $\lambda$ thermal conductivity
Mach number	$M = \frac{U}{c} = \frac{\text{speed}}{\text{sound speed}}$	(3.315)	$M$ Mach number $c$ sound speed
Rossby number	$\text{Ro} = \frac{U}{\Omega L} = \frac{\text{inertial force}}{\text{Coriolis force}}$	(3.316)	$\text{Ro}$ Rossby number $\Omega$ angular velocity

<sup>a</sup>Sometimes the square root of this expression.  $L$  is usually the fluid depth.

<sup>b</sup>Sometimes the reciprocal of this expression.

## Fluid waves

Sound waves	$v_p = \left( \frac{K}{\rho} \right)^{1/2} = \left( \frac{dp}{d\rho} \right)^{1/2}$	(3.317)	$v_p$ wave (phase) speed $K$ bulk modulus $p$ pressure $\rho$ density $\gamma$ ratio of heat capacities $R$ molar gas constant $T$ (absolute) temperature $\mu$ mean molecular mass
In an ideal gas (adiabatic conditions) <sup>a</sup>	$v_p = \left( \frac{\gamma RT}{\mu} \right)^{1/2} = \left( \frac{\gamma p}{\rho} \right)^{1/2}$	(3.318)	$v_g$ group speed of wave $h$ liquid depth $\lambda$ wavelength $k$ wavenumber $g$ gravitational acceleration $\omega$ angular frequency
Gravity waves on a liquid surface <sup>b</sup>	$\omega^2 = gk \tanh kh$	(3.319)	
	$v_g \simeq \begin{cases} \frac{1}{2} \left( \frac{g}{k} \right)^{1/2} & (h \gg \lambda) \\ (gh)^{1/2} & (h \ll \lambda) \end{cases}$	(3.320)	
Capillary waves (ripples) <sup>c</sup>	$\omega^2 = \frac{\sigma k^3}{\rho}$	(3.321)	$\sigma$ surface tension
Capillary-gravity waves ( $h \gg \lambda$ )	$\omega^2 = gk + \frac{\sigma k^3}{\rho}$	(3.322)	

<sup>a</sup>If the waves are isothermal rather than adiabatic then  $v_p = (p/\rho)^{1/2}$ .

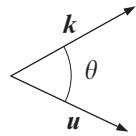
<sup>b</sup>Amplitude  $\ll$  wavelength.

<sup>c</sup>In the limit  $k^2 \gg g\rho/\sigma$ .

## Doppler effect<sup>a</sup>

Source at rest, observer moving at $u$	$\frac{v'}{v} = 1 - \frac{ \mathbf{u} }{v_p} \cos\theta$	(3.323)	$v', v''$ observed frequency $v$ emitted frequency $v_p$ wave (phase) speed in fluid
Observer at rest, source moving at $u$	$\frac{v''}{v} = \frac{1}{1 - \frac{ \mathbf{u} }{v_p} \cos\theta}$	(3.324)	$\mathbf{u}$ velocity $\theta$ angle between wavevector, $\mathbf{k}$ , and $\mathbf{u}$

<sup>a</sup>For plane waves in a stationary fluid.



## Wave speeds

Phase speed	$v_p = \frac{\omega}{k} = v\lambda$	(3.325)	$v_p$ phase speed $v$ frequency $\omega$ angular frequency ( $= 2\pi v$ ) $\lambda$ wavelength $k$ wavenumber ( $= 2\pi/\lambda$ )
Group speed	$v_g = \frac{d\omega}{dk}$	(3.326)	$v_g$ group speed
	$= v_p - \lambda \frac{dv_p}{d\lambda}$	(3.327)	

## Shocks

Mach wedge <sup>a</sup>	$\sin\theta_w = \frac{v_p}{v_b}$	(3.328)	$\theta_w$ wedge semi-angle $v_p$ wave (phase) speed $v_b$ body speed
Kelvin wedge <sup>b</sup>	$\lambda_K = \frac{4\pi v_b^2}{3g}$	(3.329)	$\lambda_K$ characteristic wavelength
	$\theta_w = \arcsin(1/3) = 19^\circ.5$	(3.330)	$g$ gravitational acceleration
Spherical adiabatic shock <sup>c</sup>	$r \simeq \left( \frac{Et^2}{\rho_0} \right)^{1/5}$	(3.331)	$r$ shock radius $E$ energy release $t$ time $\rho_0$ density of undisturbed medium
Rankine– Hugoniot shock relations <sup>d</sup>	$\frac{p_2}{p_1} = \frac{2\gamma M_1^2 - (\gamma - 1)}{\gamma + 1}$	(3.332)	1 upstream values 2 downstream values
	$\frac{v_1}{v_2} = \frac{\rho_2}{\rho_1} = \frac{\gamma + 1}{(\gamma - 1) + 2/M_1^2}$	(3.333)	$p$ pressure $v$ velocity $T$ temperature
	$\frac{T_2}{T_1} = \frac{[2\gamma M_1^2 - (\gamma - 1)][2 + (\gamma - 1)M_1^2]}{(\gamma + 1)^2 M_1^2}$	(3.334)	$\rho$ density $\gamma$ ratio of specific heats $M$ Mach number

<sup>a</sup>Approximating the wake generated by supersonic motion of a body in a nondispersive medium.

<sup>b</sup>For gravity waves, e.g., in the wake of a boat. Note that the wedge semi-angle is independent of  $v_b$ .

<sup>c</sup>Sedov–Taylor relation.

<sup>d</sup>Solutions for a steady, normal shock, in the frame moving with the shock front. If  $\gamma = 5/3$  then  $v_1/v_2 \leq 4$ .

## Surface tension

Definition	$\sigma_{lv} = \frac{\text{surface energy}}{\text{area}}$	(3.335)	$\sigma_{lv}$ surface tension (liquid/vapour interface)
	$= \frac{\text{surface tension}}{\text{length}}$	(3.336)	
Laplace's formula <sup>a</sup>	$\Delta p = \sigma_{lv} \left( \frac{1}{R_1} + \frac{1}{R_2} \right)$	(3.337)	$\Delta p$ pressure difference over surface
Capillary constant	$c_c = \left( \frac{2\sigma_{lv}}{g\rho} \right)^{1/2}$	(3.338)	$R_i$ principal radii of curvature
Capillary rise (circular tube)	$h = \frac{2\sigma_{lv} \cos \theta}{\rho g a}$	(3.339)	$c_c$ capillary constant
Contact angle	$\cos \theta = \frac{\sigma_{wv} - \sigma_{wl}}{\sigma_{lv}}$	(3.340)	$\rho$ liquid density
			$g$ gravitational acceleration
			$h$ rise height
			$\theta$ contact angle
			$a$ tube radius
			$\sigma_{wv}$ wall/vapour surface tension
			$\sigma_{wl}$ wall/liquid surface tension

<sup>a</sup>For a spherical bubble in a liquid  $\Delta p = 2\sigma_{lv}/R$ . For a soap bubble (two surfaces)  $\Delta p = 4\sigma_{lv}/R$ .

